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## *Robotics 2*

# Dynamic model of robots: Newton-Euler approach

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# Approaches to dynamic modeling

(reprise)



## energy-based approach (Euler-Lagrange)



## Newton-Euler method (balance of forces/torques)

- multi-body robot seen as a whole
  - constraint (internal) reaction forces between the links are automatically eliminated: in fact, they do not perform work
  - closed-form (symbolic) equations are directly obtained
  - best suited for study of dynamic properties and **analysis** of control schemes
- dynamic equations written separately for each link/body
  - **inverse dynamics in real time**
    - equations are evaluated in a **numeric** and **recursive** way
    - best for **synthesis** (=implementation) of model-based control schemes
  - by elimination of reaction forces and back-substitution of expressions, we still get closed-form dynamic equations (identical to those of Euler-Lagrange!)



# Derivative of a vector in a moving frame

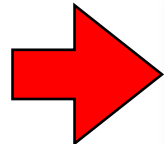
... from velocity to acceleration

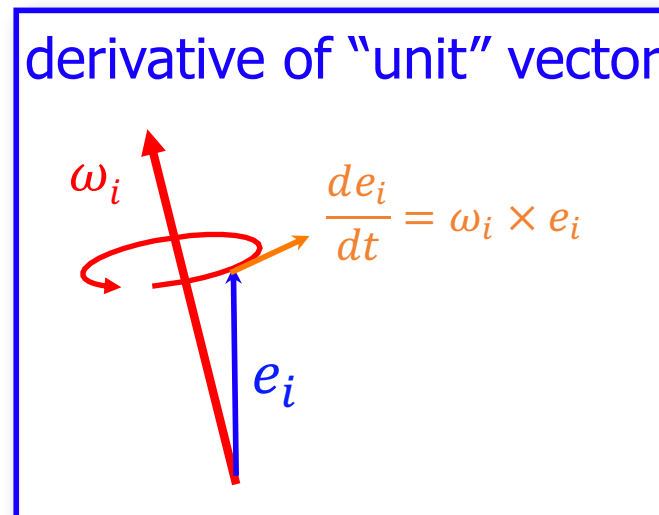
$${}^0v_i = {}^0R_i {}^i v_i$$

$${}^0\dot{R}_i = S({}^0\omega_i) {}^0R_i$$

$${}^0\dot{v}_i = {}^0a_i = {}^0R_i {}^i a_i = {}^0R_i {}^i \dot{v}_i + {}^0\dot{R}_i {}^i v_i$$

$$= {}^0R_i {}^i \dot{v}_i + {}^0\omega_i \times {}^0R_i {}^i v_i = {}^0R_i ({}^i \dot{v}_i + {}^i \omega_i \times {}^i v_i)$$


$${}^i a_i = {}^i \dot{v}_i + {}^i \omega_i \times {}^i v_i$$





# Dynamics of a rigid body

## ■ Newton dynamic equation

- **balance:** sum of forces = variation of **linear** momentum

$$\sum f_i = \frac{d}{dt}(mv_c) = m\dot{v}_c$$

## ■ Euler dynamic equation

- **balance:** sum of torques = variation of **angular** momentum

$$\begin{aligned}\sum \mu_i &= \frac{d}{dt}(I\omega) = I\dot{\omega} + \frac{d}{dt}(R\bar{I}R^T)\omega = I\dot{\omega} + (\dot{R}\bar{I}R^T + R\bar{I}\dot{R}^T)\omega \\ &= I\dot{\omega} + S(\omega)R\bar{I}R^T\omega + R\bar{I}R^T S^T(\omega)\omega = I\dot{\omega} + \omega \times I\omega\end{aligned}$$

## ■ principle of **action and reaction**

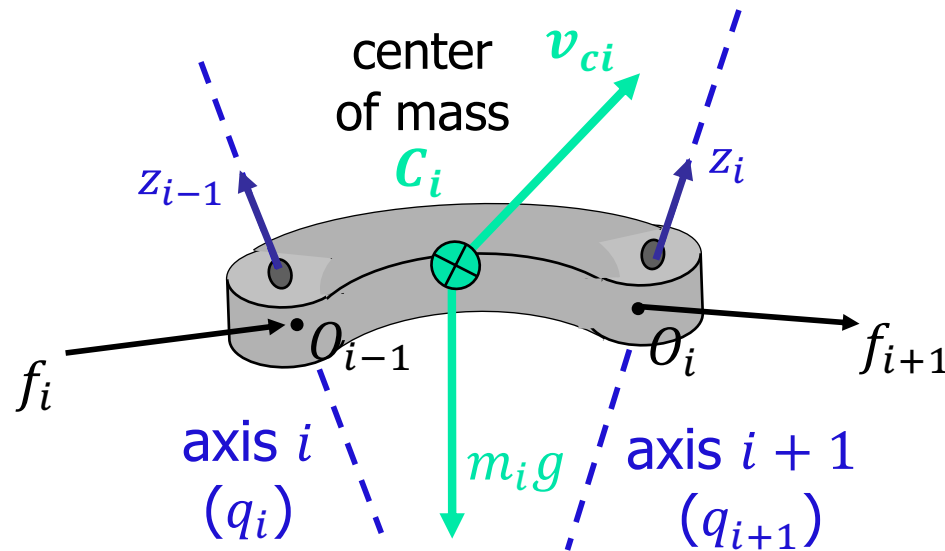
- forces/torques: applied **by** body  $i$  **to** body  $i + 1$   
= **–** applied **by** body  $i + 1$  **to** body  $i$



# Newton-Euler equations - 1

link  $i$

FORCES



$f_i$  force applied from link  $i - 1$  on link  $i$

$f_{i+1}$  force applied from link  $i$  on link  $i + 1$

$m_i g$  gravity force

all vectors expressed in the same RF (better  $RF_i$ )

Newton equation

$$f_i - f_{i+1} + m_i g = m_i a_{ci}$$

N

linear acceleration of  $C_i$



# Newton-Euler equations - 2

link  $i$

## TORQUES

$\tau_i$  torque applied from link  $(i - 1)$  on link  $i$

$\tau_{i+1}$  torque applied from link  $i$  on link  $(i + 1)$

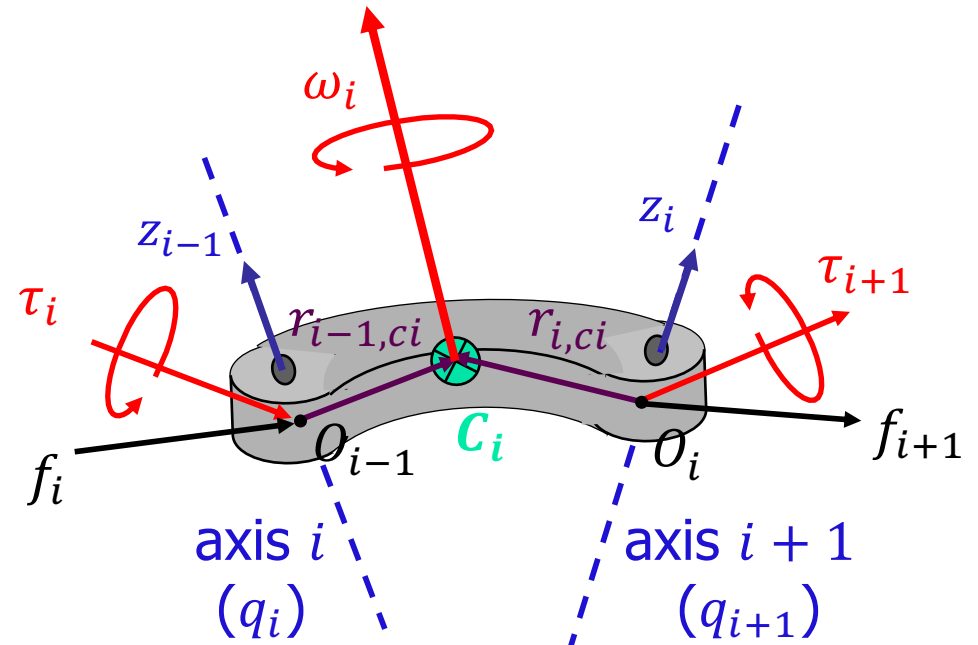
$f_i \times r_{i-1,ci}$  torque due to  $f_i$  w.r.t.  $C_i$

$-f_{i+1} \times r_{i,ci}$  torque due to  $-f_{i+1}$  w.r.t.  $C_i$

Euler equation

$$\tau_i - \tau_{i+1} + f_i \times r_{i-1,ci} - f_{i+1} \times r_{i,ci} = I_i \dot{\omega}_i + \omega_i \times (I_i \omega_i)$$

angular acceleration of body  $i$



all vectors expressed in the same RF (RF <sub>$i$</sub>  !!)

gravity force gives no torque at  $C_i$



# Forward recursion

## Computing velocities and accelerations

- “moving frames” algorithm (as for velocities in Lagrange)
- wherever there is no leading superscript, it is the same as the subscript  
( $\omega_i = {}^i\omega_i$ )
- for simplicity, only revolute joints  
(see [textbook](#) for the more general treatment)

### initializations

$$\omega_i = {}^{i-1}R_i^T [\omega_{i-1} + \dot{q}_i z_{i-1}]$$

←  $\omega_0$

$$\begin{aligned} \dot{\omega}_i &= {}^{i-1}R_i^T [\dot{\omega}_{i-1} + \ddot{q}_i z_{i-1} - \dot{q}_i z_{i-1} \times (\omega_{i-1} + \dot{q}_i z_{i-1})] \\ &= {}^{i-1}R_i^T [\dot{\omega}_{i-1} + \ddot{q}_i z_{i-1} + \dot{q}_i \omega_{i-1} \times z_{i-1}] \end{aligned}$$

AR

←  $\dot{\omega}_0$

$$a_i = {}^{i-1}R_i^T a_{i-1} + \dot{\omega}_i \times {}^i r_{i-1,i} + \omega_i \times (\omega_i \times {}^i r_{i-1,i})$$

←  $a_0 - {}^0g$

$$a_{ci} = a_i + \dot{\omega}_i \times r_{i,ci} + \omega_i \times (\omega_i \times r_{i,ci})$$

the gravity force term can be skipped in Newton equation, if added here



# Backward recursion

## Computing forces and torques

from  $N_i$   $\longrightarrow$  to  $N_{i-1}$ 
eliminated, if inserted  
in forward recursion ( $i=0$ )
initializations

$f_i = f_{i+1} + m_i(a_{ci} - \cancel{g})$ 
 $\longleftarrow f_{N+1}$ 
 $\tau_{N+1}$

$\tau_i = \tau_{i+1} - f_i \times (r_{i-1,i} + r_{i,c_i}) + f_{i+1} \times r_{i,c_i} + I_i \dot{\omega}_i + \omega_i \times (I_i \omega_i)$

from  $E_i$   $\longrightarrow$  to  $E_{i-1}$

F/TR

at each step of this recursion, we have two vector equations ( $N_i + E_i$ ) at the joint providing  $f_i$  and  $\tau_i$ : these contain ALSO the **reaction forces/torques** at the joint axis  $\Rightarrow$  they should be "projected" next along/around this axis

FP
 $u_i = \begin{cases} f_i^T \cdot z_{i-1} + \eta_i \dot{q}_i & \text{for prismatic joint} \\ \tau_i^T \cdot z_{i-1} + \eta_i \dot{q}_i & \text{for revolute joint} \end{cases}$ 
 $\longrightarrow$ 

 $N$  scalar equations at the end

generalized forces  
(in rhs of Euler-Lagrange eqs)

add here dissipative terms  
(here viscous friction only)





# Comments on Newton-Euler method

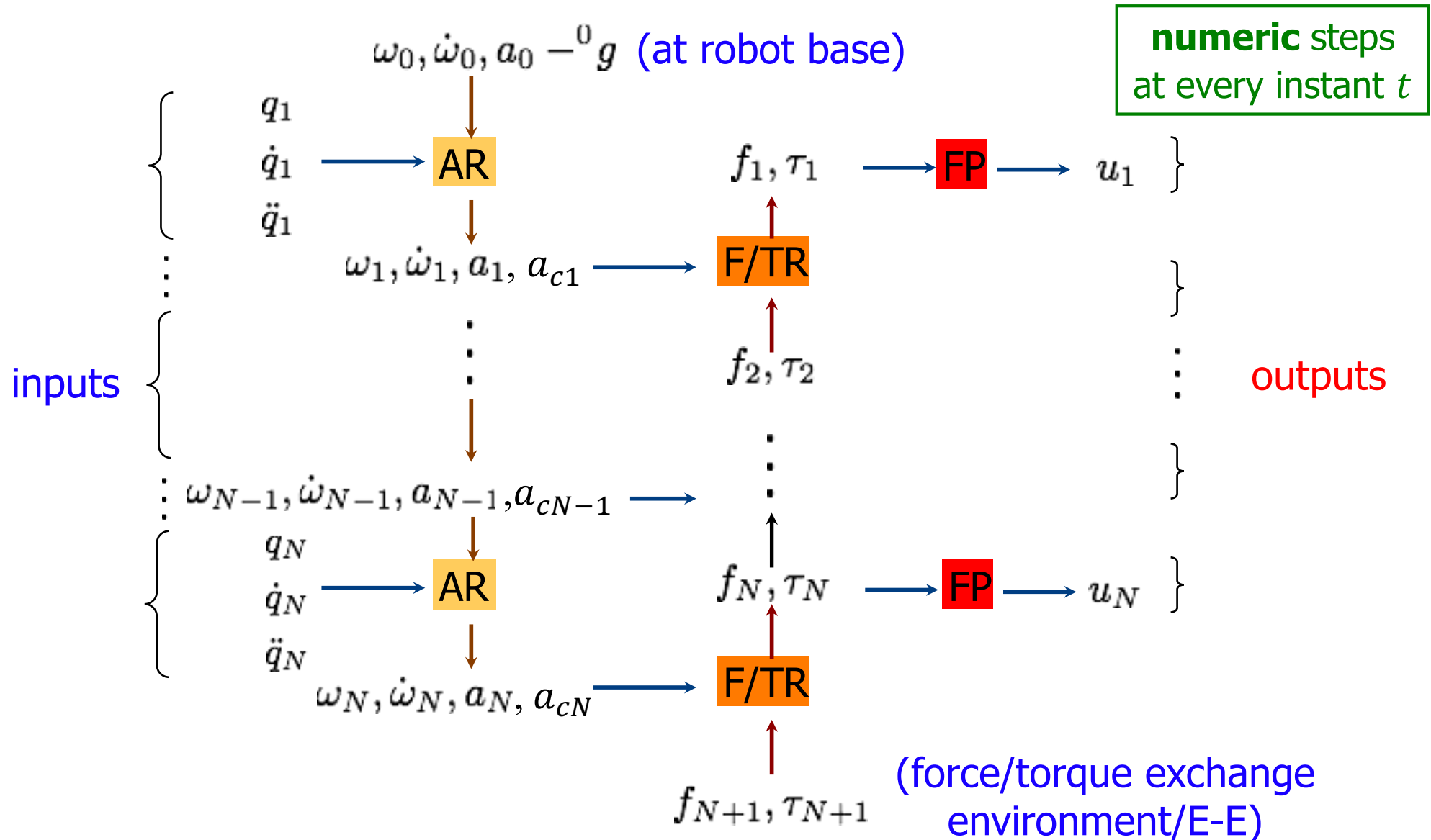
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- the previous forward/backward recursive formulas can be evaluated in symbolic or numeric form
  - **symbolic**
    - substituting expressions in a recursive way
    - at the end, a closed-form dynamic model is obtained, which is identical to the one obtained using Euler-Lagrange (or any other) method
    - there is **no** special convenience in using N-E in this way
  - **numeric**
    - substituting numeric values (numbers!) at each step
    - **computational complexity** of each step remains constant  $\Rightarrow$  grows **in a linear fashion** with the number  $N$  of joints ( $O(N)$ )
    - strongly recommended for real-time use, especially when the number  $N$  of joints **is large**



# Newton-Euler algorithm

efficient computational scheme for inverse dynamics





# Matlab (or C) script

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general routine  $NE_\alpha(\text{arg}_1, \text{arg}_2, \text{arg}_3)$

- data file (of a specific robot)
  - number  $N$  and types  $\sigma = \{0,1\}^N$  of joints (revolute/prismatic)
  - table of DH kinematic parameters
  - list of **ALL** dynamic parameters of the links (and of the motors)
- input
  - vector parameter  $\alpha = \{^0g, 0\}$  (presence or absence of gravity)
  - three ordered **vector arguments**
    - typically, samples of joint **position, velocity, acceleration** taken from a desired trajectory
- output
  - generalized force  $u$  for the **complete** inverse dynamics
  - ... or **single terms** of the dynamic model



# Examples of output

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- complete inverse dynamics

$$u = NE_0g(q_d, \dot{q}_d, \ddot{q}_d) = M(q_d)\ddot{q}_d + c(q_d, \dot{q}_d) + g(q_d) = u_d$$

- gravity terms

$$u = NE_0g(q, 0, 0) = g(q)$$

- centrifugal and Coriolis terms

$$u = NE_0(q, \dot{q}, 0) = c(q, \dot{q})$$

- $i$ -th column of the inertia matrix

$$u = NE_0(q, 0, e_i) = M_i(q)$$

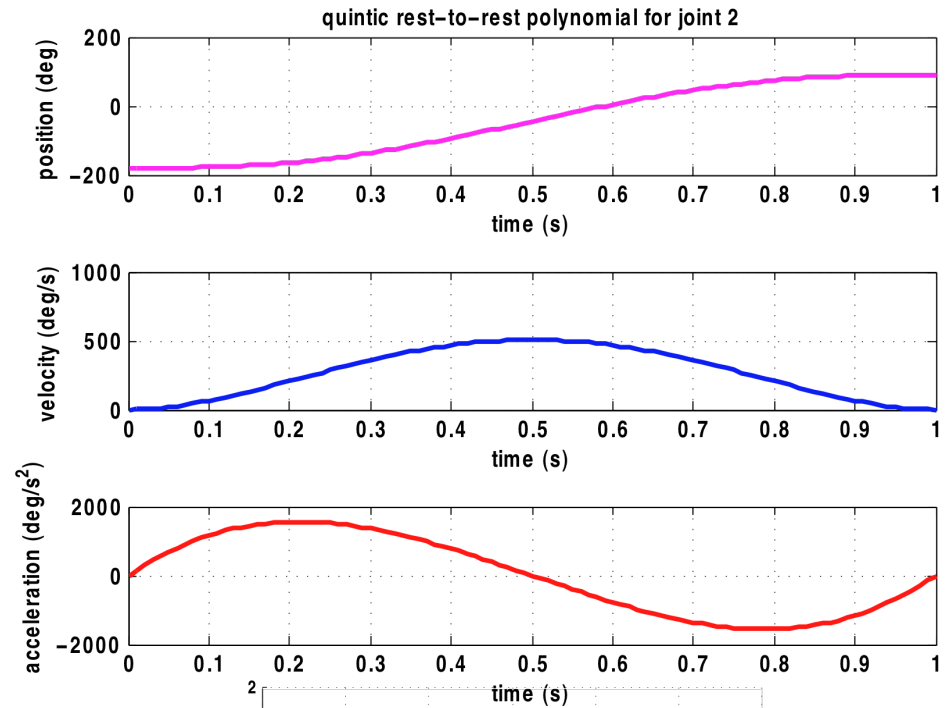
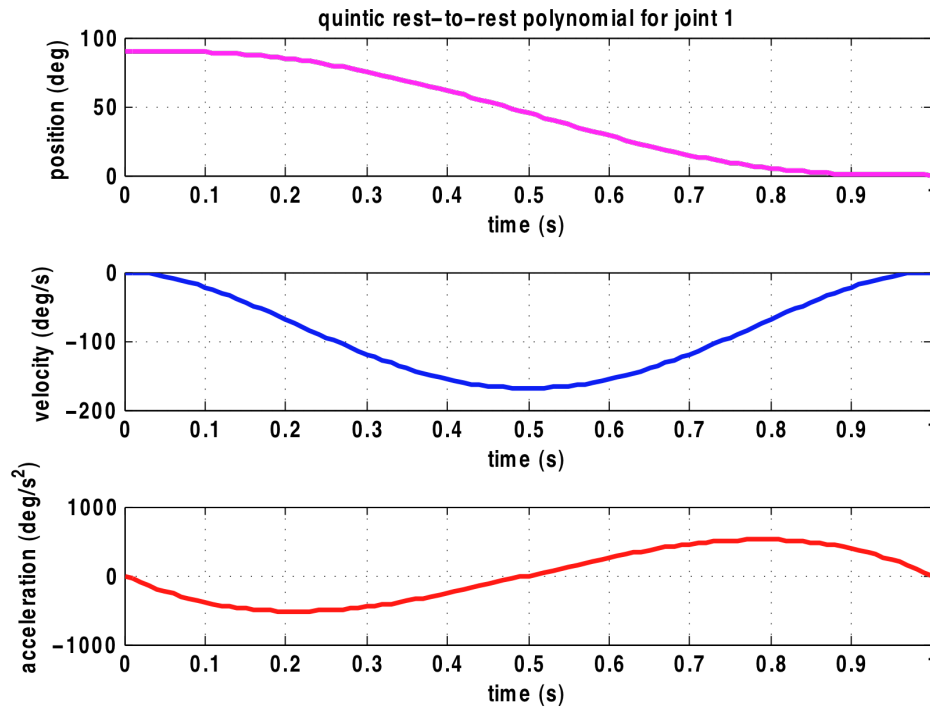
$e_i = i$ -th column  
of identity matrix

- generalized momentum

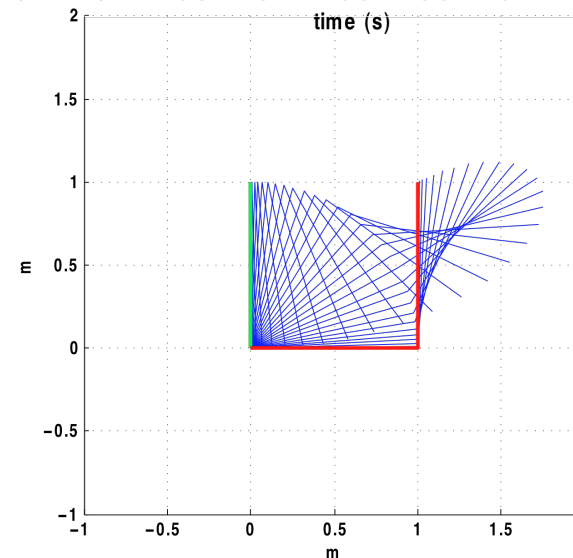
$$u = NE_0(q, 0, \dot{q}) = M(q)\dot{q} = p$$



# Inverse dynamics of a 2R planar robot

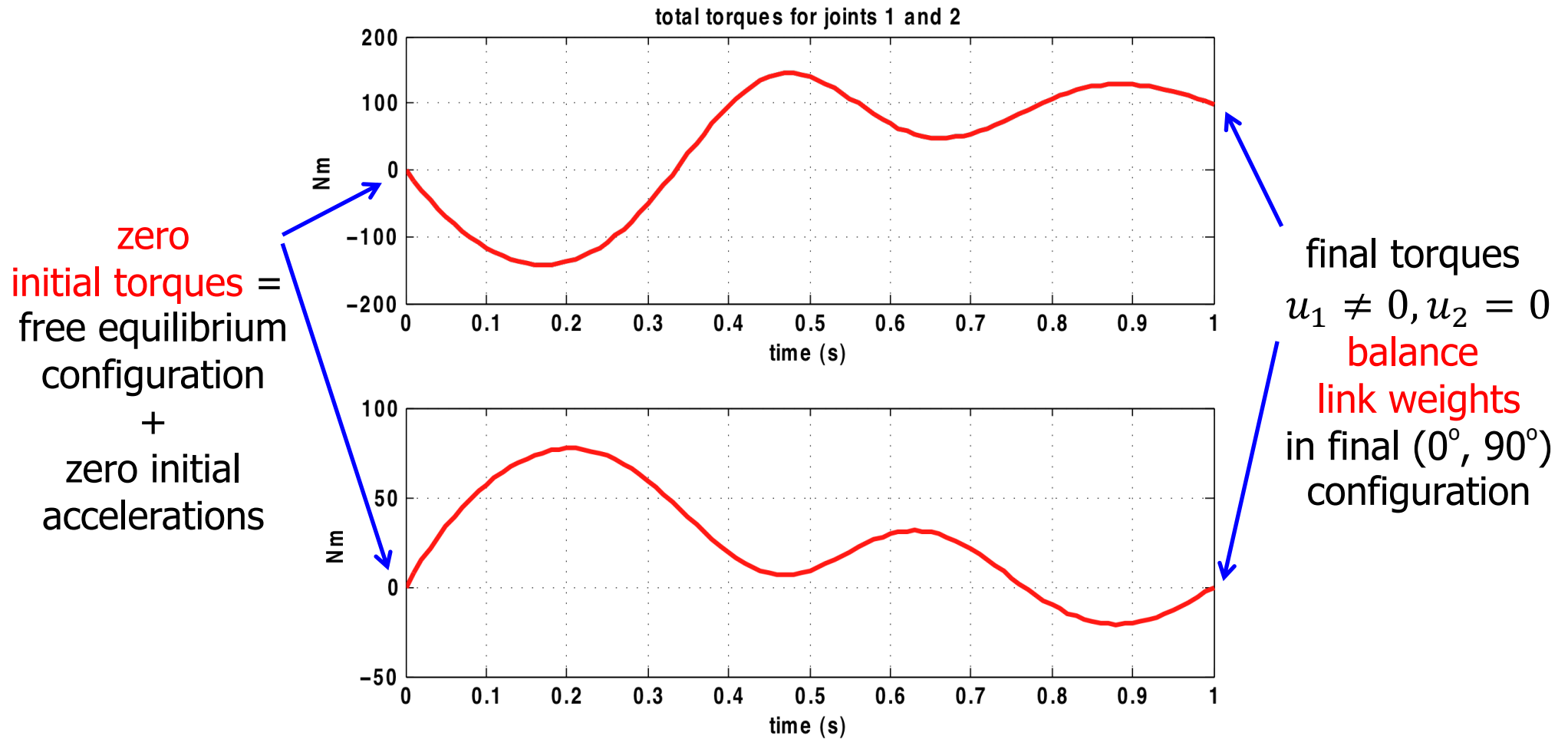


desired (smooth) joint motion:  
quintic polynomials for  $q_1, q_2$  with  
zero vel/acc boundary conditions  
from  $(90^\circ, -180^\circ)$  to  $(0^\circ, 90^\circ)$  in  $T = 1$  s



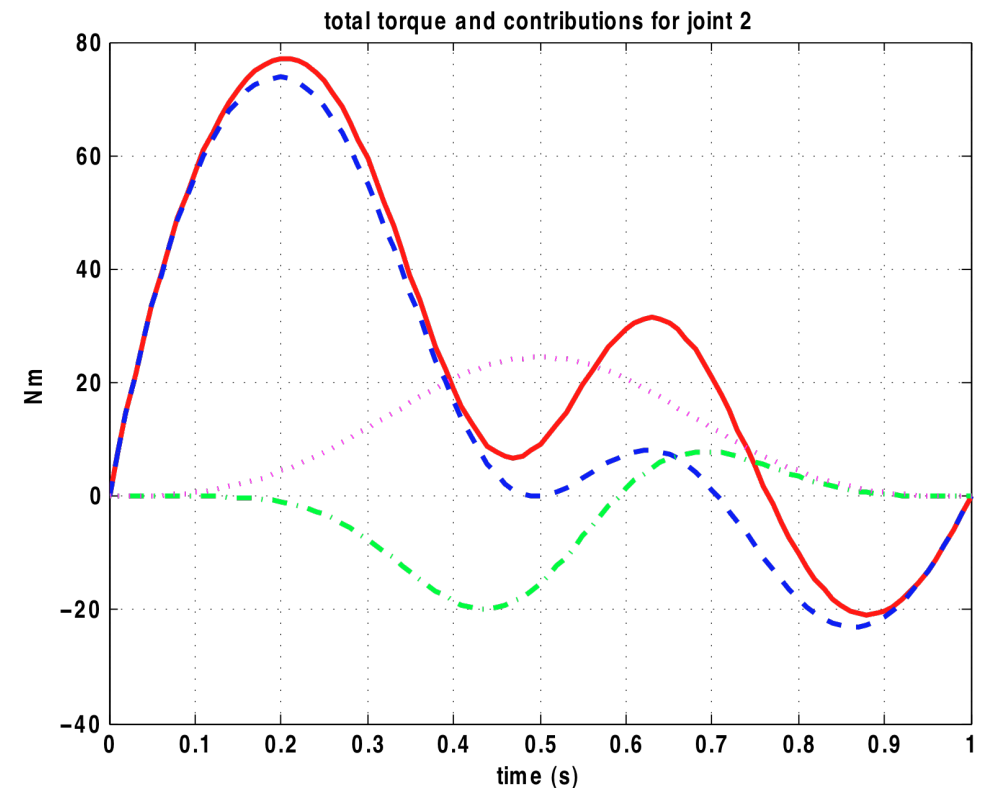
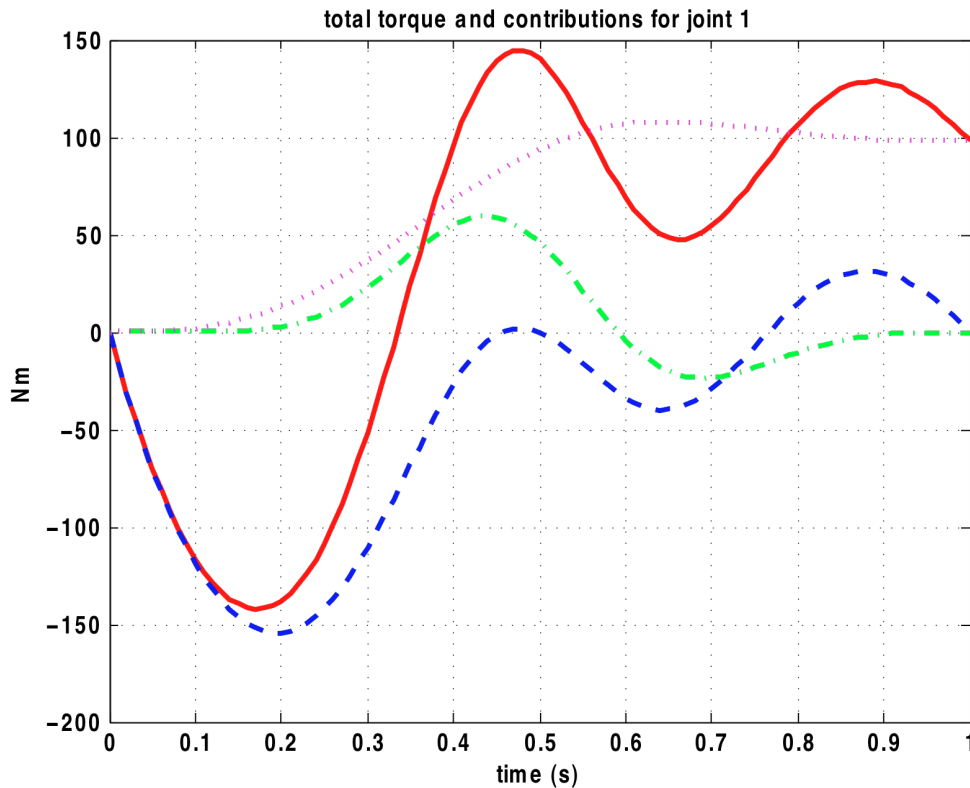


# Inverse dynamics of a 2R planar robot



motion in vertical plane (under gravity)  
both links are thin rods of uniform mass  $m_1 = 10$  kg,  $m_2 = 5$  kg

# Inverse dynamics of a 2R planar robot



torque contributions at the two joints for the desired motion

— = total, --- = inertial  
- · - · - = Coriolis/centrifugal, ····· = gravitational

# Use of NE routine for simulation

## direct dynamics



- numerical integration, at **current** state  $(q, \dot{q})$ , of  
$$\ddot{q} = M^{-1}(q)[u - (c(q, \dot{q}) + g(q))] = M^{-1}(q)[u - n(q, \dot{q})]$$
- Coriolis, centrifugal, and gravity terms

$$n = NE_0(q, \dot{q}, 0) \quad \text{complexity } O(N)$$

- $i$ -th column of the inertia matrix, for  $i = 1, \dots, N$

$$M_i = NE_0(q, 0, e_i) \quad O(N^2)$$

- numerical inversion of inertia matrix

$$InvM = \text{inv}(M) \quad O(N^3) \\ \text{but with small coefficient}$$

- given  $u$ , integrate acceleration computed as

$$\ddot{q} = InvM * [u - n] \quad \longrightarrow \quad \text{new state } (q, \dot{q}) \\ \text{and repeat over time ...}$$